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On the Simplifications of Cable Model in Static Analysis of Large-Dimension Cable-Driven Parallel Robots

Dinh Quan Nguyen, Marc Gouttefarde, Olivier Company and François Pierrot

Abstract—This paper addresses the simplification of cable model in static analysis of large-dimension cable-driven parallel robots (CDPR). An approach to derive a simplified hefty cable model is presented. The approach provides an insight into the limitation of such a simplification. The resulting cable tension computation is then used to solve the inverse kinematic problem of CDPR. A new expression of cable length taking into account both the non-negligible cable mass and elasticity is also introduced. Finally, simulations and experiments on a large CDPR prototype are provided. The results show that taking into account both cable mass and elasticity improves the robot accuracy.

I. INTRODUCTION

In the past twenty years, cable-driven parallel robots (CDPR) have been extensively studied in favor of their appealing advantages such as light-weight, simple in design and relatively low cost. Compared to parallel manipulators with rigid links, CDPR use cables to position their end-effector directly by controlling the length of each cable. Possible applications in manufacturing industry, construction industry and aerospace industry are positioning and handling of large and heavy parts across wide workspaces. Several studies on large-dimension CDPR have been made [1]–[15].

In many previous studies on CDPR, all cables have been considered massless. This assumption is usually valid for robots of reasonable size and carrying light payloads, but for large-dimension robots or for robots that carry heavy payloads, such an assumption may be invalid [13]. Irvine presents in [16] the well-known cable model known as the elastic catenary which takes the cable mass and elasticity into account. In [17], a cable catenary model in 3D Cartesian space is introduced. Because of their highly nonlinear nature, numerical methods must be considered to deal with the catenary cable modeling. However, these methods are usually relatively slow because of the involved computations. Therefore, using such a cable model for CDPR may be impractical in real-time environments, especially in solving the inverse kinematics and tension distribution problems. Finding means of reducing the complexity of such cable models is thus of great practical interest.

Over the past decades, the methods to deal with the catenary cable modeling have been simplified and improved. In [16], a parabolic cable model is presented. It is valid if the sagging of the cable is small enough. However, the suggested validity condition is not explicitly derived. In [7], Kozak presents a method to solve the inverse kinematics of general CDPR when the cable model is the elastic catenary. An expression for the unstrained cable length, which takes into account the cable mass but without considering the cable elasticity, is provided. Rui Yao et al. introduce in [12] the large-dimension four-cable-driven parallel robot in FAST in which the parabolic cable model known in [16] is used to obtain an expression for the cable length. However, the latter is just an approximation of the one presented in [7] and does not account for the cable elasticity. In [13], a simplified static analysis of large-dimension CDPR is introduced based on the parabolic hefty cable model presented in [16]. The determination of the cable tension distribution is thereby transformed into one that is similar in form to the case of CDPR with massless cables. The domain of validity of the analysis proposed in [13] was left as an open issue. Later, in [14], Hui Li et al. give a relationship between the cable horizontal and vertical force components which neglects the elasticity effect from the catenary cable model. No additional development of the cable model from this relationship are proposed by these authors.

Efficient means to find practical forms of cable model have been proposed. Nevertheless unsolved issues still exist. Some of them are the validity of the simplified static analysis proposed in [13] and the precision of the inverse kinematic solutions given in [7]. The motivation of the present paper comes directly from these issues. First of all, a rigorous simplification procedure of cable catenary model in static analysis is presented. It provides a linear relationship between the cable force components and gives a clear view of the involved simplifications as well as the limitation of the method. Then, a solution for the computation of cable length is demonstrated. To the best of our knowledge, the resulting expression of the cable length is new. It takes into account both the non-negligible mass and elasticity of the cable.

The paper is organized as follows. The well-known cable catenary is recalled in Section II. The rigorous simplification steps are introduced in Section III and leads to the same parabolic simplified cable model as the one presented in [16]. This model leads directly to a linear relationship between cable horizontal and vertical force components. The limitation of the simplified cable model is discussed in Section IV. The new expression of the cable length taking into account both the cable mass and elasticity is derived in Section V. Finally, Section VI presents simulation and experimental results obtained for a large-dimension 6-DOF.
parallel robot driven by \( m = 8 \) cables.

II. CABLE PROFILE

Let us consider a steel cable that has unstrained length \( L_0 \ (m) \), self-weight \( w \ (N/m) \), elastic modulus \( E \ (Pa) \) and cable cross-section area \( A_0 \ (m^2) \).

Fig. 1 shows the relevant coordinates and parameters of a cable lying in a vertical plane in static equilibrium. The cable is fixed between two end-points \( A(x_a, z_a) \) and \( B(x_b, z_b) \). All the coordinates are in the local frame attached to the vertical plane containing the cable. The term \( \Delta L \) represents the strain of the cable.

A point \( M \) along the strained cable has Cartesian coordinates \( x \) and \( z \). The variable \( p \) represents the strained length of the cable segment as measured from the end-point \( A \) of the cable to the point \( M \). The variable \( s \) will be used to denote the unstrained length of the same cable segment. In Fig.1, the origin of the local frame will be placed at the end-point \( A \). Thus, the variable \( s \) lies in the range:

\[
0 \leq s \leq L_0
\]

with \( M(s = 0) = A \) and \( M(s = L_0) = B \).

The well-known catenary equations can be written as follows [16]:

\[
x(s) = x_b + \frac{\tau_{bx}(s - L_0)}{E A_0} + \frac{\tau_{bx}}{w} \ln \left[ \frac{\tau_a + \tau_{bz} + w(s - L_0)}{\tau_b + \tau_{bz}} \right] \tag{1a}
\]

\[
z(s) = z_b + \frac{\tau_{bz}(s - L_0)}{E A_0} + \frac{w(s - L_0)^2}{2EA_0} + \frac{1}{w}(\tau_s - \tau_b) \tag{1b}
\]

where \( \tau_{bx}, \tau_{bz} \) are the cable horizontal and vertical force components at point \( B \) and \( \tau_s, \tau_b \) are the tensions in cable at point \( M \) and \( B \), respectively:

\[
\tau_b = \sqrt{\tau_{bx}^2 + \tau_{bz}^2} \tag{3}
\]

The shape of the cable must satisfy the geometric constraint:

\[
\left( \frac{dx}{dp} \right)^2 + \left( \frac{dz}{dp} \right)^2 = 1 \tag{4}
\]

which implies:

\[
\frac{dp}{dx} = \sqrt{1 + \left( \frac{dz}{dx} \right)^2} \tag{5}
\]

At point \( M \), the force balance for the segment of the cable between points \( M \) and \( B \) can be written as follows:

\[
\tau_s \left( \frac{dx}{dp} \right) = \tau_{bx} \tag{6a}
\]

\[
\tau_s \left( \frac{dz}{dp} \right) = \tau_{bz} + w(s - L_0) \tag{6b}
\]

The cable tension at point \( M \) is considered to satisfy Hooke’s law:

\[
\tau_s = E A_0 \left( \frac{dp}{ds} - 1 \right) \tag{7}
\]

The relationship between coordinates \( x \) and \( z \) which does not depend explicitly on the variable \( s \) can be derived from (5), (6) and (7):

\[
\frac{d^2z}{dx^2} = \frac{w}{\tau_{bx}} \cdot \frac{E A_0}{\frac{E A_0 + \tau_s}{\tau_{bx}} \cdot \sqrt{1 + \left( \frac{dz}{dx} \right)^2}} \tag{8}
\]

Let us assume that \( \tau_s \ll E A_0 \). Then, the cable elasticity has a very little influence on the cable shape since (8) can be reduced to the simpler expression:

\[
\frac{d^2z}{dx^2} = \frac{w}{\tau_{bx}} \cdot \sqrt{1 + \left( \frac{dz}{dx} \right)^2} \tag{9}
\]

The solution of (9) is:

\[
z(x) = \frac{\tau_{bx}}{w} \cosh \left( \frac{w}{\tau_{bx}} x + C_1 \right) + C_2 \tag{10}
\]

which must satisfy the following boundary conditions:

\[
z(x_a) = z_a \tag{11a}
\]

\[
z(x_b) = z_b \tag{11b}
\]

As the origin of the local cable frame has been chosen at the end-point \( A \), \( x_a = 0 \) and \( z_a = 0 \), then:

\[
C_2 = -\frac{1}{\mu} \cosh (C_1) \tag{12}
\]

\[
z_b = \frac{1}{\mu} \left[ \cosh (\mu x_b + C_1) - \cosh (C_1) \right] \tag{13}
\]

where

\[
\mu = \frac{w}{\tau_{bx}} \tag{14}
\]

The term \( C_1 \) can be found by solving (13):

\[
C_1 = \ln \left( \frac{\sqrt{\mu^2 z_b^2 + e^{\mu x_b} + e^{-\mu x_b} - 2 + \mu z_b}}{e^{\mu x_b} - 1} \right) \tag{15}
\]
The condition to achieve (15) are:
\[ x_b > 0 \quad (16a) \]
\[ \tau_{bx} > 0 \quad (16b) \]

(16a) is verified by definition of the local cable frame. (16b) is always true because the cable can only pull on the mobile platform. Note that (15) is the only solution of (13) because conditions (16a) and (16b) are always verified. The tangent at point B of the cable is then computed as:
\[ \tan(\beta) = \frac{dz}{dx} \bigg|_{x=x_b} = \frac{\tau_{bz}}{\tau_{bx}} = \sinh(\mu x_b + C_1) \quad (17) \]

Relationship (17) is similar to the one presented in [14]. Note that the simplifying assumption \( EA_0 \gg \tau_s \) and the validity conditions (16a) and (16b) are not explicit in [14].

III. LINEARIZATION

In the case of cable-driven parallel robots with massless cables, several methods have been developed to solve important issues such as the tension distribution problem for 6-DOF CDPR with a number of cables \( m \geq 6 \), e.g. in [19] and [20]. To take advantage of such methods, as presented in [13], in case of CDPR with cables of non-negligible mass, one must achieve a linear relationship between the horizontal and vertical cable force components in order to transform the equilibrium equations of the mobile platform into ones similar to the case of massless CDPR. From (17), since \( \tan(\beta) \) is a function of the cable horizontal force component \( \tau_{bx} \), the relationship between the cable force components at end-point B is a nonlinear relationship \( \tau_{bz} = \tan(\beta) \cdot \tau_{bx} \). Hence, simplification of the term \( \tan(\beta) \) needs to be performed in a proper way. To this end, one can consider the Taylor series expansion of (17) in terms of the variable \( \mu \) defined in (14).

In fact, choosing \( \mu \) as the expansion variable is reasonable since, in the case of a CDPR with steel cables carrying heavy payloads, \( \mu \) is relatively small and depends only on the cable self-weight \( w \) and the cable horizontal force component \( \tau_{bx} \). The Taylor series expansion of \( \tan(\beta) \) around the expansion point \( \mu = 0 \) up to order of 3 is:
\[ \tan(\beta) = -\sinh(H) + \frac{1}{2} x_b \cosh(H) \mu \]
\[ - \frac{x_b^2 \sinh(H)}{24 \sqrt{x_b^2 + z_b^2}} + \frac{1}{8} x_b^2 \sinh(H) \mu^2 \]
\[ + \frac{1}{48} x_b^3 \cosh(H) + \frac{x_b^2 \sinh(H)}{48 \sqrt{x_b^2 + z_b^2}} \] \( \mu^3 + O(\mu^4) \) \quad (18)

where
\[ H = \ln \left( \frac{x_b}{\sqrt{x_b^2 + z_b^2 + z_b^2}} \right) \]

In order to achieve a linear relationship between the cable force components \( \tau_{bx} \) and \( \tau_{bz} \), only expansion of order 1 in (18) is to be considered:
\[ \tan(\beta) = -\sinh(H) + \frac{1}{2} x_b \cosh(H) \mu \]
\[ = \frac{z_b}{x_b} + \frac{1}{2} \sqrt{x_b^2 + z_b^2} \mu \quad (19) \]

Finally, we achieve the linear relationship:
\[ \tau_{bz} = \tan(\beta_0) \cdot \tau_{bx} + \frac{wL}{2} \quad (20) \]

where \( \tan(\beta_0) = z_b/x_b \) and \( L = \|AB\| = \sqrt{x_b^2 + z_b^2} \).

Relationship (20) is the same as the simplified cable model presented in [13].

In summary, we have rigorously detailed the simplifications to the elastic catenary cable modeling which lead to a linear relationship between the horizontal cable force \( \tau_{cx} \) and vertical cable force components in (20). This linear relationship is the basis of the simplified static analysis introduced in [13]. The first assumption \( \tau_s \ll EA_0 \) was made in (8). If verified, it means that the cable elasticity has no real influence on the cable shape. The second assumption states that the variable \( \mu \) defined in (14) is small enough. Depending on the context, the validity of these two assumptions may be an issue so that they should be carefully checked.

IV. LIMITATION OF THE SIMPLIFIED CABLE MODEL

In the linearization procedure presented in Section III, if the expansion variable is chosen as \( \nu = \mu x_b \) instead of \( \mu \), and the Taylor series expansion of \( \tan(\beta) \) is performed around the expansion point \( \nu = 0 \), we obtain the same linear relationship (20). Hence, the accuracy of the approximation of \( \tan(\beta) \) should also be justified in term of the variable \( \nu \) since \( \nu > \mu \) for large-dimension CDPR \( (x_b \text{ can be large}) \). Equation (20) is valid if \( \nu \) is small enough:
\[ 0 \leq \nu \leq \nu_{\text{min}} \quad (21) \]

where \( \nu_{\text{min}} \) is a chosen small value. This condition leads to the following constraint on the cable horizontal force component:
\[ \frac{w x_b}{\nu_{\text{min}}} \leq \tau_{bx} \quad (22) \]

The validation of condition (22) gives an insight into the limitation of the simplified cable model. To illustrate this point, let us consider the problem which consists in determining the wrench-feasible workspace (WFW) [18] of a CDPR driven by \( m \) cables. The WFW is the set of poses of the mobile platform for which, for any wrench \( \mathbf{f} \) in \( W_{\text{req}} \), there exists a vector of cable horizontal force components \( \tau_{bx} \) which is a solution to the equilibrium equations:
\[ \mathbf{W} \tau_{bx} = \mathbf{f} \quad (23) \]

and verifies
\[ \tau_{\text{min}} \leq \tau_{bx,i} \leq \tau_{\text{max}} \quad (i = 1, m) \quad (24) \]

Here, \( W_{\text{req}} \) is the required set of wrenches that the cables must apply at the reference point \( P_i \) of the mobile platform. Equation (23) is the modified equilibrium equations of the mobile platform obtained with the simplified static analysis presented in [13] which assumes that (20) is valid. For a given CDPR geometry, the wrench matrix \( \mathbf{W} \) depends only on the pose of the mobile platform. If the required set of
wrenches $W_{req}$ is also known, then the size of the WFW is determined by condition (24).

For the simplified static analysis of [13] to be valid across the WFW, the lower bound of (24) must not be smaller than

$$\tau_{\text{min}} = \frac{w x_b}{\nu_{\text{min}}} \quad (25)$$

The size of the WFW is proportional to the admissible range $[\tau_{\text{min}}, \tau_{\text{max}}]$ of the cable horizontal force component. In fact, the upper bound on the cable forces is always determined by safety considerations. Hence, with a fixed upper bound on the cable horizontal force component ($\tau_{\text{max}} = \text{const}$), the size of the part of the WFW in which the simplified cable model is valid is reduced when $\tau_{\text{min}}$ increases. That makes the size of this workspace inversely proportional to $x_b$ (and also to the cable self-weight $w$).

In practice, it is preferable to have a large WFW. For a given set of wrenches and in a large desired WFW, condition (22) may be violated in some areas. Thus, it is expected to get large errors in the approximation of the term $\tan(\beta)$. Such errors lead to a poor accuracy in solving the cable tension distribution problem with the simplified static analysis of [13]. The problem is even more involved since the assumption $\tau_s \ll EA$ may be questioned if $\tau_b$ is large.

V. COMPUTATION OF THE CABLE LENGTH

Starting from (7), one can derive:

$$d\tau = \left(1 + \frac{\tau_s}{EA_0}\right) ds = \left(1 + \frac{\tau_{bx}}{EA_0} \cdot \frac{dp}{dx}\right) ds$$

$$\frac{dp}{dx} = \left(1 + \frac{\tau_{bx}}{EA_0} \frac{dp}{ds} \frac{ds}{dx} \right) \quad (26)$$

Substitute $dp/dx$ from (5) into (26) to obtain:

$$\frac{ds}{dx} = \frac{1}{1 + \frac{\tau_{bx}}{EA_0} \left(1 + \left(\frac{dz}{dx}\right)^2\right)^2} \quad (27)$$

If we use the solution of $z(x)$ from (10) and since,

$$\sqrt{1 + \left(\frac{dz}{dx}\right)^2} = \cosh (\mu x + C_1) \quad (28)$$

where $C_1$ is given in (15), then (27) becomes:

$$\frac{ds}{dx} = \frac{EA_0}{\tau_{bx}} \left(1 - \frac{1}{1 + R}\right) \quad (29)$$

where

$$R = \frac{\tau_{bx}}{EA_0} \cosh (\mu x + C_1)$$

Since $R \ll 1$ because we assume that $\tau_{bx} \ll EA_0$, let us use the Taylor series expansion to find an approximation of $ds/dx$ around the expansion point $R = 0$ as:

$$\frac{ds}{dx} = \frac{EA_0}{\tau_{bx}} \left[ R - R^2 + R^3 + O(R^4) \right] \quad (30)$$

The cable unstrained length can finally be obtained as follows:

$$L_0 = \int_0^{L_0} ds = \int_{x=0}^{x_b} \frac{EA_0}{\tau_{bx}} \left[ R - R^2 + R^3 \right] dx$$

$$= \frac{1}{\mu} \left[ \sinh(C_0) - \sinh(C_1) \right]$$

$$- \frac{\tau_{bx}}{4E A_0 \mu} \left[ 2\mu x_b - \sinh(2C_1) + \sinh(2C_b) \right]$$

$$- \frac{\tau_{bx}^3}{12 E^2 A_0^2 \mu} \left[ 9 \sinh(C_1) - 9 \sinh(C_b) + \sinh(3C_1) - \sinh(3C_b) \right] \quad (31)$$

with $C_b = \mu x_b + C_1$.

Let us note that, in (31), if the series expansion of $ds/dx$ only up to order 1 of $R$ is considered, the expression of $L_0$ is the same as the one presented in [7] where only the cable mass is taken into account (elasticity neglected).

The unstrained cable length as given in (31) is the one needed for the inverse kinematics and thus for the trajectory planning part in a control scheme of a CDPR.

VI. VALIDATION

The validation of the simplified cable model with the new cable length expression (31) is performed in simulation and experimentally on the CoGiRo prototype, a large-dimension suspended CDPR driven by $m = 8$ cables shown in Fig. 2. The characteristics of the steel cables driving the mobile platform are:

- $E = 35e9$ (Pa)
- $A_0 = 8.2051e-06$ (m$^2$)
- $w = 0.62784$ (N/m)

A. Simulation

The center of mass $P_c$ of the mobile platform of the robot follows a desired trajectory in the Cartesian space $X_t = (x_t \ y_t \ z_t)$ with constant orientation $\alpha_{x,y,z} = (0 \ 0 \ 0)$ (X-Y-Z Euler angle convention). The mass of the mobile platform is $m_p = 500$ kg. $X_t$ is the position of $P_c$ at time $t$. Smoothed trapezoidal velocity method [22] is used to generate the desired trajectory. The robot is simulated
using the massless cable model \( L_{0i\text{Massless}} = \| A_i B_i \| \), the elastic cable model (where only elasticity is considered) and the simplified hefty cable model with different expressions of cable unstrained length corresponding to Taylor series expansion of \( L_0 \) up to order 1 and 2 of \( R \) in (31). The cable unstrained lengths of the elastic cable model are computed as follows:

\[
L_{0i\text{Elastic}} = \frac{L_{0i\text{Massless}}}{1 + \frac{\tau_i}{EA_0}} \quad (i = 1, m)
\]

where \( \tau_i \) is the cable tension of the \( i \)-th cable obtained in the case of massless cable model.

Different criteria are used to evaluate the results:

\[
dL = \max_{1 \leq i \leq m} |L_{0i} - L_{0i\text{Catenary}}|
\]

\[
dX_{\text{Model}} = \| X_{t\text{Model}} - X_t \|
\]

where \( L_{0i} \) is the unstrained length of the \( i \)-th cable in the cases of the massless, elastic and simplified hefty cable models, \( L_{0i\text{Catenary}} \) is the unstrained length of the \( i \)-th cable in case of the elastic catenary cable model. \( X_{t\text{Model}} \) is the simulated position of \( P_i \) at time \( t \) obtained for the CDPR with a given cable model.

The variable \( \nu \) (introduced in Section IV) is computed in case of the simplified hefty cable models to illustrate the discussion of Section IV:

\[
\nu_{\text{max}} = \max_{1 \leq i \leq m} \frac{w_{2bi}}{\tau_{bri}}
\]

To check the assumption \( \tau_s \ll EA_0 \) (see Section III), we compute also the ratio

\[
\eta_{\text{max}} = \max_{1 \leq i \leq m} \frac{\tau_i}{EA_0}
\]

The method presented in [20] is used to solve the tension distribution problem in case of massless cable model and simplified static analysis [13]. The results are then compared to the one obtained by means of the elastic catenary cable model which is considered to be the reference.

Figure 3 shows that hefty cable models lead to significant improvements compared to the massless cable model in solving the inverse kinematics (unstrained cable length). The massless cable model and simplified hefty cable model which has expression of \( L_{0i} \) with order 1 of \( R \) in (31) (\( \text{Simplified}_{R1} \) in Fig. 3) have a similar accuracy. The simplified cable model which uses the expression of \( L_{0i} \) with order 2 of \( R \) in (31) (considering both cable mass and cable elasticity, \( \text{Simplified}_{R2} \) in Fig. 3) gives the best accuracy. Note that, in term of solving the inverse kinematics, elastic cable model and \( \text{Simplified}_{R2} \) give similar results. It means that the elasticity has a large impact on the cable lengths. This is reasonable for the steel cables with the given characteristics (the term \( EA_0 \) is relatively small).

In the second test, the center of mass of the mobile platform of the robot is fixed at \( P_c = (0 \ 0 \ 0) \) in the Cartesian space. The payload of the mobile platform is varied in the range \( 100 \leq m_p \leq 800 \) kg. Figure 4 shows that the cable model \( \text{Simplified}_{R2} \) gives the best results. The impact of elasticity becomes more visible when the payload increases.

Now, let us consider the two assumptions suggested in
Section III for the simplified cable models to be valid. The
evolution of the two variables $\nu_{\text{max}}$ and $\eta_{\text{max}}$ in Fig. 4 shows
a tradeoff of the simplified models. When the payload is
small, $\eta_{\text{max}}$ is small, thus the first assumption $\tau_s \ll EA_0$ is
valid but $\nu_{\text{max}}$ is large, hence the second assumption $\nu \ll 1$
is violated and vice versa. According to this example, it
can be seen that the validity of simplified cable model is a
difficult issue when solving kinetostatic problems of CDPR.

B. Experiments

In our tests, the sagging effect of the cables is only
considered between the end-points $A_1$ and $B_1$. The cable
segments from end-points $A_i$ to the winch drums are ver-
tical (no sagging) and assumed unstressed (the deformation
effects are compensated by adding some offsets to the cable
lengths). The total mass of the mobile platform in our tests
is $m_y = 210\text{ kg}$. The position of the center of mass $P_c$
of the mobile platform is measured by a laser tracker. The
cable characteristics are given in Section VI-A. The tension
distribution problem is solved by using the method in [20]
based on static analysis of [13].

The robot CoGiRo is controlled by means of the scheme
presented in [21]. In the latter, inverse kinematics is used to
compute the desired cable unstrained lengths. To this end,
(31) has been used in the experiments reported in the present
paper.

The robot positioning accuracy is evaluated at several
static equilibrium poses inside its workspace. Figure 5 shows
the positions of these points (top view). The criteria (34),
(35) and (36) are used to verify the results. In Tab. 1,
$X_L$, $X_E$, $X_{R1}$, $X_{R2}$ are the measured positions of the
center of mass of the mobile platform in case of the massless,
estatic, $\text{Simplified}_{R1}$ and $\text{Simplified}_{R2}$ cable models, respec-
tively.

Table I shows that the cable model $\text{Simplified}_{R2}$ gives
the best results. The elastic cable model has accuracies
quite close to $\text{Simplified}_{R2}$. The massless cable model and
$\text{Simplified}_{R1}$ provide similar results with a lower accuracy.

This is reasonable since we achieved similar performances
in simulation for the inverse kinematics (Fig. 3, ”Differences
in cable unstrained length”). In the case of the CDPR at
hand (CoGiRo prototype), the differences in experimental
results confirm the strong impact of the cable elasticity on
the robot accuracy ($dX_E$ in Tab. 1), notably as compared to
the cable mass alone ($dX_{R1}$). It remains that the best results
are obtained by considering both cable mass and elasticity
($dX_{R2}$).

In Tab. 1, at equilibrium points that are close to the
center of the practical workspace ($\nu_{\text{max}}$ is small), we achieve
good accuracies (e.g. at points $P_1$ and $P_2$). The accuracy of
the robot reduces when the mobile platform is positioned
near the workspace boundary (e.g. at points $P_4 - P_8$).
Indeed, the value of the variable $\nu_{\text{max}}$ increases when $P_c$
approaches the workspace boundary (when $x_c$ and $y_c$ are
large) which creates expected errors in solving the inverse
kinematic problem. In addition, the variable $\eta_{\text{max}}$ having
relatively large values, it also reduces the accuracy of the
simplified cable models. This fact confirms our analysis of
the limitation of the simplified cable model in Section IV.

CONCLUSIONS

A simplification of the catenary cable model has been
discussed in this paper. The resulting simplified cable model
Corresponds to the parabolic hefty cable model presented
in previous studies. However, in these studies, most of the
assumptions for this simplified cable model to be valid
were not justified. Our approach provides a clear view of
the limitation of such a simplified cable model which was
verified in simulation and experiments. Indeed, the analysis
of the examples in Section VI illustrated that the conditions
for the simplified cable model to be valid are not always true
and, hence, that they should be checked carefully.

Furthermore, a new expression of the cable unstrained
length has been introduced. It accounts for both the cable
mass and longitudinal elasticity. This new expression of
the cable unstrained length greatly helps to improve the
positioning accuracy of a large-dimension CDPR as shown
by the experimental results presented in this paper. To fur-
ther improve positioning accuracy, the cable characteristics,
notably its Young’s modulus $E$, should be experimentally
identified which is part of our ongoing works.

Although this paper focus mainly on static analysis, the
expression of the cable unstrained length introduced here
can be useful in a dynamic analysis where the dynamics of
the mobile platform and the actuators that drive the cables
are taken into account. The dynamics of the cables can be
neglected especially when heavy payloads are handled due to
slow motions of the mobile platform. Further studies about
the impacts of the presented simplified cable models on
dynamic performances will be discussed in our future works.

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TABLE I: Positioning accuracy of the center of mass of the mobile platform of CoGiRo at several static equilibrium poses

<table>
<thead>
<tr>
<th>Static poses</th>
<th>Xc</th>
<th>Yc</th>
<th>Zc</th>
</tr>
</thead>
<tbody>
<tr>
<td>(x_c, y_c, z_c) (m)</td>
<td>(x_c, y_c, z_c) (m)</td>
<td>(x_c, y_c, z_c) (m)</td>
<td>(x_c, y_c, z_c) (m)</td>
</tr>
<tr>
<td>P1</td>
<td>(0.0, 0.0, 1.4)</td>
<td>(-0.001, 0.000, 0.822)</td>
<td>(-0.001, 0.000, 0.822)</td>
</tr>
<tr>
<td>P2</td>
<td>(0.0, 0.0, 3.0)</td>
<td>(-0.001, 1.999, 1.041)</td>
<td>(-0.001, 1.999, 1.041)</td>
</tr>
<tr>
<td>P3</td>
<td>(-0.0, 4.0, 0.2)</td>
<td>(-2.997, 1.999, 0.144)</td>
<td>(-2.998, 1.999, 0.144)</td>
</tr>
<tr>
<td>P4</td>
<td>(0.0, 0.3, 0.5)</td>
<td>(4.999, 3.999, 0.444)</td>
<td>(4.999, 3.999, 0.444)</td>
</tr>
<tr>
<td>P5</td>
<td>(4.0, 3.0, 1.0)</td>
<td>(4.000, -1.999, 0.911)</td>
<td>(4.000, -1.999, 0.911)</td>
</tr>
</tbody>
</table>

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REFERENCES